

المحاضرة الخامسة

Chapter (2)

The relation between field and circuit theory: Maxwell's equations.

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2.1 Introduction :

In circuit theory we deal with circuit element , the voltage (V) across them , and the total current (I) through them . In field theory we deal with the field vectors (E,D,B,H,and J)and their values as a function of position .

Consider a short rod of length (l) and cross-sectional area(A) , Thus the voltage difference between the ends of the rod is , from ohms law ,

$$V= IR \dots\dots\dots (1)$$

Where (I) is the current through the rod

From ohms law at point

$$E= \frac{J}{\sigma} \dots\dots\dots(2)$$

Where J =conduction current density, Am⁻²

σ = conductivity, Ωm^{-1}

Now, integrating (2) over the length of the rod, we obtain the voltage difference (V) between the ends. That is

$$V = \int E \cdot dl = \int \frac{J}{\sigma} \cdot dl \dots\dots(3)$$

$$= \frac{Jl}{\sigma} = JA \frac{l}{\sigma A} \dots\dots(4)$$

Where JA=I = current through rod ,A

$l/\sigma A=R$ = resistance of rod , Ω

A=cross section area of rod , m²

Thus from (4) we have

$$V=IR \dots\dots(5)$$

Starting with field theory , we have arrived at the circuit relation known as ohm's law .

2.2 Maxwell's equation as generalizations of circuit equations

$$\nabla \times H = J + \frac{\partial D}{\partial t} \dots\dots\dots(1)$$

Equation (1) it is the differential form of Maxwell's equation as derived from Ampers law

$$\nabla \times E = - \frac{\partial B}{\partial t} \dots\dots\dots(2)$$

Equation (2) it is the differential form of Maxwell's equation as derived from faradays law

$$\nabla \cdot D = \rho \dots\dots\dots(3)$$

Equation (3) is Maxwell's electric field equation as derived from Gauss law in differential form

$$\nabla \cdot B = 0 \dots\dots\dots(4)$$

Equation (4) is Maxwell's magnetic field equation as derived from Gauss's law

2-3 Maxwell's equation in free space

Maxwell's equation are stated in their general form , for the special case of free space , where the current density (J) and the charge density (ρ) are zero , the equations reduce to a simple form – In integral form the equations are

$$\oint H \cdot dl = \int_s \frac{\partial D}{\partial t} \cdot ds \dots\dots\dots(1)$$

$$\oint E \cdot dl = - \int_s \frac{\partial B}{\partial t} \cdot ds \dots\dots\dots(2)$$

$$\oint_s D \cdot ds = 0 \dots\dots\dots(3)$$

$$\oint_s B \cdot ds = 0 \dots\dots\dots(4)$$

-In Differential form the equations are

$$\nabla \times H = \frac{\partial D}{\partial t} \dots\dots\dots(5)$$

$$\nabla \times E = - \frac{\partial B}{\partial t} \dots\dots\dots(6)$$

$$\nabla \cdot D = 0 \dots \dots \dots (7)$$

$$\nabla \cdot B = 0 \dots \dots \dots (8)$$

2.4 Maxwell's equation for harmonically varying fields

If we assume that the fields vary harmonically with time, Maxwell's equations can be expressed in another special form.

Thus, if (D) varies with time as given by

$$D = D_0 e^{j\omega t} \dots \dots \dots (1)$$

Then $\frac{\partial D}{\partial t} = j\omega D_0 e^{j\omega t} = j\omega D \dots \dots (2)$

When the same assumption is made for B,

$$B = B_0 e^{j\omega t}$$

$$\frac{\partial B}{\partial t} = j\omega B_0 e^{j\omega t} = j\omega B, \quad B = \mu H$$

$$\frac{\partial B}{\partial t} = j\omega \mu H \dots \dots \dots (3)$$

Maxwell's equations in integral form reduce to

$$\oint H \cdot dl = (\sigma + j\omega\epsilon) \int_s E \cdot ds \dots \dots (4)$$

$$\oint E \cdot dl = -j\omega\mu \int_s H \cdot ds \dots \dots (5)$$

$$\oint_s D \cdot ds = \int_s \rho dV \dots \dots (6)$$

$$\oint_s B \cdot ds = 0 \dots \dots (7)$$

In differential form they are

$$\nabla \times H = (\sigma + j\omega\epsilon)E \dots \dots (8)$$

$$\nabla \times E = -j\omega\mu H \dots\dots\dots(9)$$

$$\nabla \cdot D = \rho \dots\dots\dots(10)$$

$$\nabla \cdot B = 0 \dots\dots\dots(11)$$

2.5 Electromagnetic wave in vacuum

The wave equation for E and B

In regions of space, there is no charge or current, Maxwell's equations read

$$\nabla \cdot E = 0$$

$$\nabla \cdot B = 0$$

$$\nabla \times E = - \frac{\partial B}{\partial t}$$

$$\nabla \times B = \mu_0 \epsilon_0 \frac{\partial E}{\partial t}$$

Electromagnetic wave equation for the field vector \vec{E} By applying the curl to equation(3) we get

$$\nabla \times E = - \frac{\partial B}{\partial t} \rightarrow \nabla \times \nabla \times E = \nabla \times \left(- \frac{\partial B}{\partial t} \right)$$

$$\rightarrow \nabla \times \nabla \times E = - (\nabla \times B)$$

By using the identity : $\nabla \times \nabla \times E = \nabla (\nabla \cdot E) - \nabla^2 E$.

$$\nabla (\nabla \cdot E) - \nabla^2 E = - \frac{\partial}{\partial t} (\nabla \times B) \rightarrow \text{but } \nabla \cdot E = 0 \text{ from (1)}$$

$$\text{and } \nabla \times B = \mu_0 \epsilon_0 \frac{\partial E}{\partial t}$$

$$\nabla^2 E = \mu_0 \epsilon_0 \frac{\partial^2 E}{\partial t^2} \dots\dots\dots(1)$$

but $\epsilon_0\mu_0 = 1/c^2$

The relation (1) becomes $\nabla^2 E = 1/c^2 \frac{\partial^2 y}{\partial t^2}$ (2)

Electromagnetic wave equation for the field vector \vec{B} By applying the curl to equation (4) we get

$$\begin{aligned} \nabla \times B &= \mu_0\epsilon_0 \frac{\partial E}{\partial t} \longrightarrow \nabla \times \nabla \times B = \nabla \times (\mu_0\epsilon_0) \frac{\partial E}{\partial t} \\ &\longrightarrow \nabla \times \nabla \times B = \mu_0\epsilon_0 \frac{\partial}{\partial t} (\nabla \times E) \end{aligned}$$

By using the identity $\nabla \times \nabla \times B = \nabla (\nabla \cdot B) - \nabla^2 B$.

$$\nabla (\nabla \cdot B) - \nabla^2 B = \mu_0\epsilon_0 \frac{\partial}{\partial t} (\nabla \times E) \longrightarrow \text{but } \nabla \cdot B = 0 \quad \text{from (2)}$$

$$\text{and } \nabla \times E = - \frac{\partial B}{\partial t} \quad \text{from (3)}$$

$$-\nabla^2 B = - \mu_0\epsilon_0 \frac{\partial}{\partial t} \left(\frac{\partial B}{\partial t} \right) \quad \nabla^2 B = \mu_0\epsilon_0 \frac{\partial^2 B}{\partial t^2} \quad \text{.....(3)}$$

but $\epsilon_0\mu_0 = \frac{1}{c^2}$

the relation between the electric \vec{E} and the magnetic vector \vec{B} is

$$E = cB \quad \text{.....(5)}$$

Where (C) is the speed of light or the speed of electromagnetic wave and equal $C = 3 \times 10^8$ m/s.

EX / A parallel – plate capacitor of radius (R) and separation (d) has a voltage at the center given

by $V = V_0 \sin \omega t$. as function of radius (r) (for $r < R$) find (a) the displacement density $J_d(r)$

(b) the magnetic field $H(r)$, take $d \ll R$

$$\text{Sol} / \oint_c H \cdot dl = \int_s (J_c + J_d) \cdot ds$$

$$\oint_c H \cdot dl = \int_s (\sigma E + \frac{\partial D}{\partial t}) \cdot ds \longrightarrow J_c = \sigma E = 0 \longrightarrow \sigma = 0$$

$$\text{a/ } J_d = \int_s (\epsilon_0 \frac{\partial E}{\partial t}) \cdot ds$$

$$J_d = \epsilon_0 \frac{d}{dt} \left(\frac{v_0 \sin \omega t}{d} \right) \longrightarrow J_d = \frac{\epsilon_0 \omega v_0 \cos \omega t}{d}$$

$$\text{b/ } \oint_c H \cdot dl = \int_s J_d \cdot ds$$

$$H \cdot 2\pi r = \frac{\epsilon_0 \omega v_0 \cos \omega t}{d} \cdot \pi r^2$$

$$H = \frac{\epsilon_0 \omega v_0 \cos \omega t \cdot \pi r^2}{2d \cdot \pi r}$$

$$H = \frac{r \epsilon_0 v_0 \cos \omega t}{2d}$$