

University of Anbar
College of Science
Department of Physics



فيزياء الحالة الصلبة Solid state Physics

المرحلة الرابعة
الكورس الاول
صلبة 1

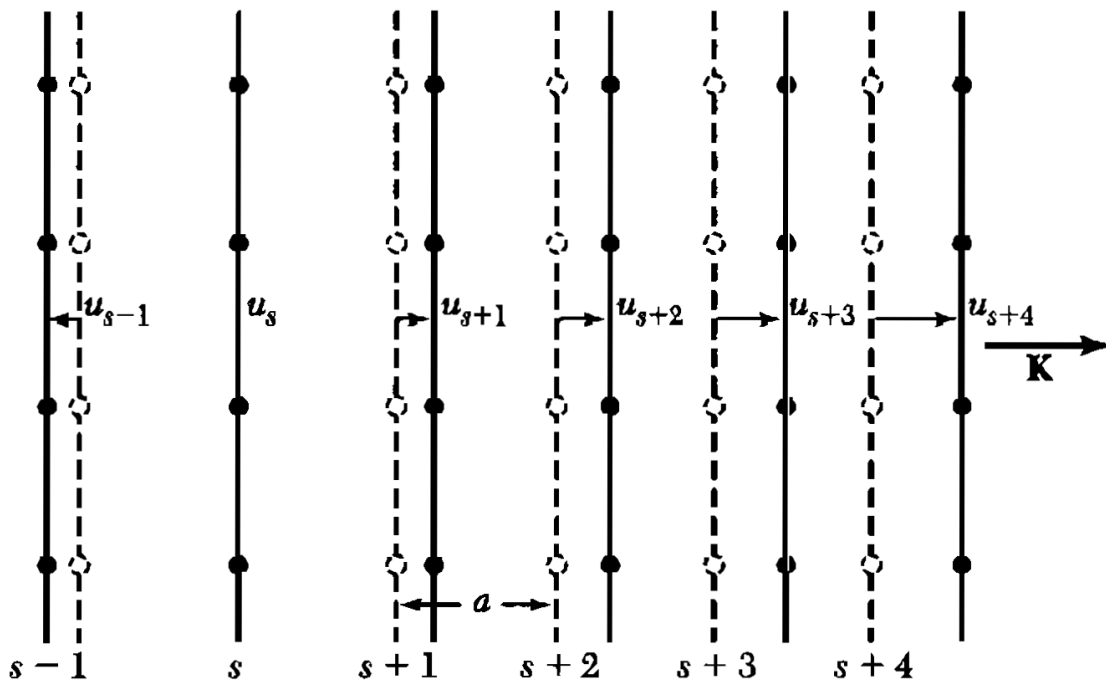
Syllabus of this course مفردات الكورس

- 1- crystal structure
- 2- x-ray diffraction
- 3- lattice vibration
- 4- thermal properties in solids
- 5- free electron model

4. Lattice or Crystal Vibrations اهتزاز الشبكة او البلورة

4.1 Vibrations Of Crystals With Monatomic احادية الذرة Basis

- Consider **نعتبر** the elastic vibrations **الاهتزاز مرنة** of a crystal with one atom in the primitive cell. We want to find the frequency of an elastic wave in terms of the wavevector that describes **لوصف** the wave **الموجة** and in terms of the elastic constants.
- Assume **نفترض** that the elastic response **الاستجابة مرنة** of the crystal is a linear function **دالة خطية** of the forces.
- We assume that the force on the planes caused by the displacement **الازاحة** of the plane $s + p$ is proportional **تناسب** to the difference $u_{s+1} - u_s$ of their displacements.



- we consider only nearest-neighbor **التفاعل** interactions **الجار الاقرب**, with $p \pm 1$, the total force on s from plane $s \pm 1$.

$$F_s = C(u_{s+1} - u_s) + C(u_{s-1} - u_s) . \quad (1)$$

This expression **التعبير** is linear **خطية** in the displacements **الازاحات** and is of the form **قانون هوك** of Hook's law **وبصيغة**.

The constant C is the force constant between nearest-neighbor planes.

The equation of motion of an atom in the plane s is

$$M \frac{d^2 u_s}{dt^2} = C(u_{s+1} + u_{s-1} - 2u_s) , \quad (2)$$

where M is the mass of an atom.

We look for solutions with all displacements having the time dependence $\exp(-i\omega t)$. then $d^2 u_s / dt^2 = -\omega^2 u_s$, and equation (2) becomes

$$-M\omega^2 u_s = C(u_{s+1} + u_{s-1} - 2u_s) . \quad (3)$$

This is a difference equation in the displacements u and has traveling wave solutions of the form:

$$u_{s\pm 1} = u \exp(isKa) \exp(\pm iKa) , \quad (4)$$

where a is the spacing between planes and K is the wavevector. The value to use for a will depend on the direction of K .

From equations 3, 4 we have

$$-\omega^2 M u \exp(isKa) = Cu[\exp[i(s+1)Ka] + \exp[i(s-1)Ka] - 2 \exp(isKa)] \quad (5)$$

We cancel $u \exp(isKa)$ from both sides, to leave

$$\omega^2 M = -C[\exp(iKa) + \exp(-iKa) - 2] \quad (6)$$

With the identity $2 \cos Ka = \exp(iKa) + \exp(-iKa)$, we have the **dispersion relation** $\omega(K)$.

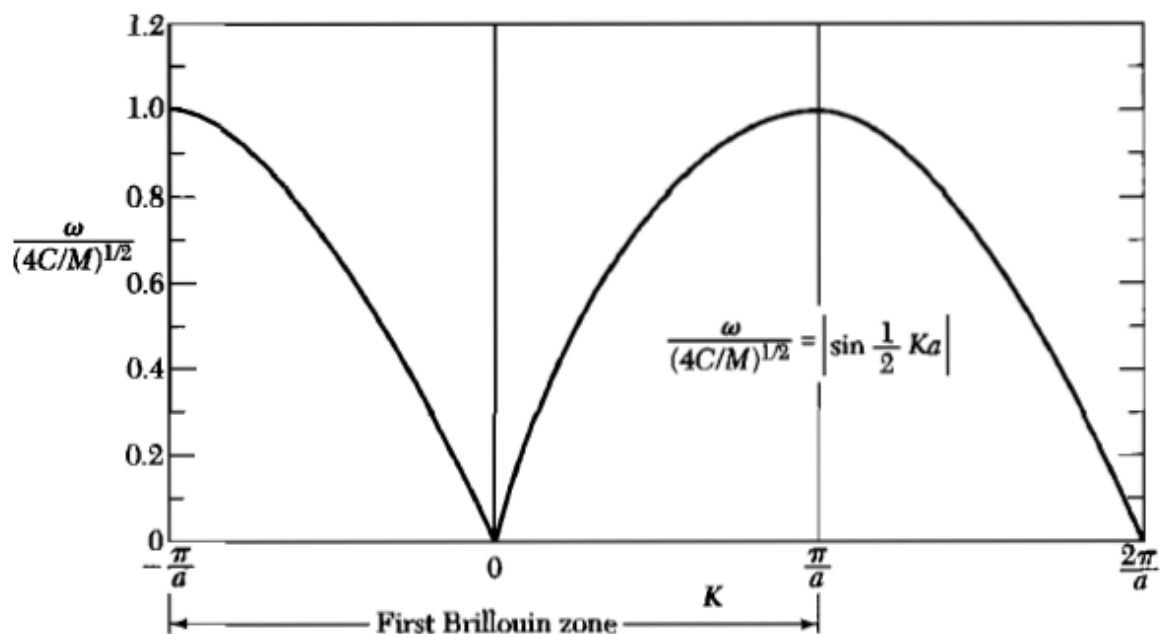
$$\omega^2 = (2C/M)(1 - \cos Ka) \quad (7)$$

The boundary of the first Brillouin zone lies at $K = \pm\pi/a$. We show from (7) that the slope of ω versus K is zero at the zone boundary:

$$d\omega^2/dK = (2Ca/M) \sin Ka = 0 \quad (8)$$

at $K = \pm\pi/a$, for here $\sin Ka = \sin(\pm\pi) = 0$.

$$-\pi < Ka \leq \pi, \quad \text{or} \quad -\frac{\pi}{a} < K \leq \frac{\pi}{a}$$



Group velocity:

The transmission velocity of a wave packet is the **group velocity**, given as

$$v_g = d\omega/dK$$

4.2 Vibrations Of Crystals With Two Atoms per primitive Basis

We consider a cubic crystal where atoms of mass M_1 lie on one set of planes and atoms of mass M_2 , lie on planes interleaved between those of the first

$$\begin{aligned} M_1 \frac{d^2 u_s}{dt^2} &= C(v_s + v_{s-1} - 2u_s) ; \\ M_2 \frac{d^2 v_s}{dt^2} &= C(u_{s+1} + u_s - 2v_s) . \end{aligned} \quad (10)$$

The solutions of these equations are:

$$u_s = u \exp(isKa) \exp(-i\omega t) ; \quad v_s = v \exp(isKa) \exp(-i\omega t) . \quad (11)$$

From equations 10, 11

$$-\omega^2 M_1 u = Cv[1 + \exp(-iKa)] - 2Cu ;$$

$$-\omega^2 M_2 v = Cu[\exp(iKa) + 1] - 2Cv .$$

$$\begin{vmatrix} 2C - M_1 \omega^2 & -C[1 + \exp(iKa)] \\ -C[1 + \exp(iKa)] & 2C - M_2 \omega^2 \end{vmatrix} = 0 ,$$

or

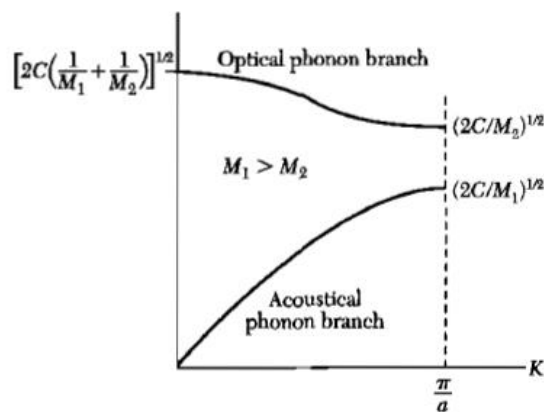
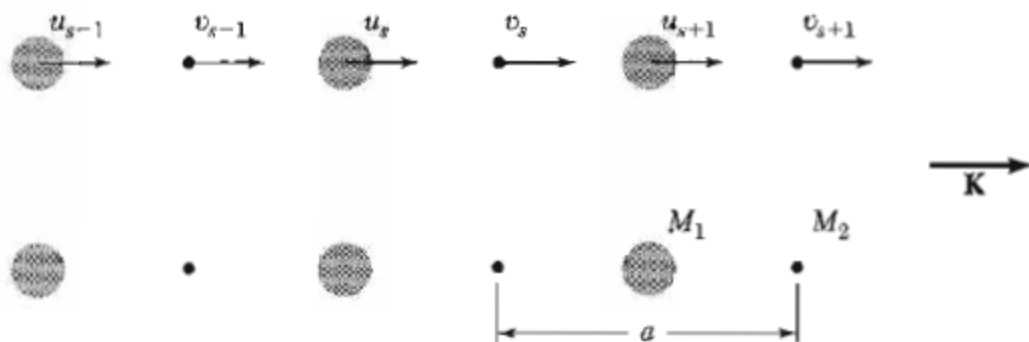
$$M_1 M_2 \omega^4 - 2C(M_1 + M_2)\omega^2 + 2C^2(1 - \cos Ka) = 0 .$$

$$\omega^2 \cong 2C \left(\frac{1}{M_1} + \frac{1}{M_2} \right) \quad (\text{optical branch}) ;$$

$$\omega^2 \cong \frac{\frac{1}{2}C}{M_1 + M_2} K^2 a^2 \quad (\text{acoustical branch}) .$$

The extent of the first Brillouin zone is $-\pi/a \leq K \leq \pi/a$, where a is the repeat distance of the lattice. At $K_{\text{max}} = \pm\pi/a$ the roots are

$$\omega^2 = 2C/M_1 ; \quad \omega^2 = 2C/M_2 .$$



References

- 1- Charles Kittel - Introduction to Solid State Physics-Wiley (2005)
- 2- J. S. Blakemore - Solid State Physics-Cambridge University Press (1985)
- 3- M. A. OMAR Elementary-solid-state-physics